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ION ACOUSTIC TURBULENCE, ANOMALOUS TRANSPORT, AND SYSTEM DYNAMICS IN HALL EFFECT THRUSTERS

Robert Martin¹, Jonathan Tran²

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²ERC INC.,
EDWARDS AIR FORCE BASE, CA USA



IPAM Mathematics of Turbulence Retreat, June 2017
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U.S. AIR FORCE



1 INTRODUCTION

2 TRANSPORT

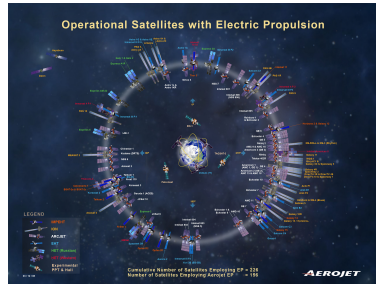
3 DYNAMIC SYSTEM

4 SUMMARY AND CONCLUSION



EP-Devices:

- EP Improves Thrust/Mass (Isp)
- Ion/Elec.-Thermal/HET Flying
- FRC/MPD/Electrospray/etc. in Dev.
- Space notoriously Risk Adverse (i.e. Tech must be “Proven”)



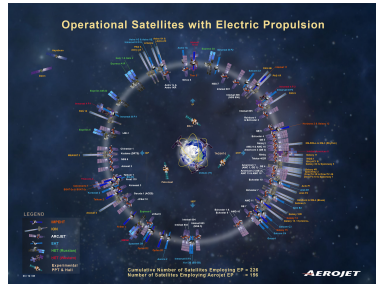
Aerojet Overview of EP Satellites (3/08)



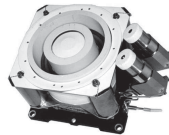
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- Plasma still Threat Thruster/Plume
- Improved Life Estimates Needed
- Space not Replicated in Ground Test

Models needed to Bridge Gap



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Fakel SPT-100



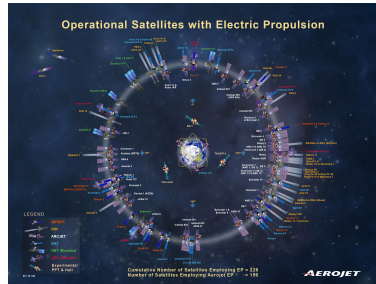
NASA Hermes Thruster



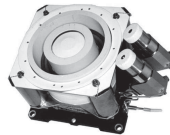
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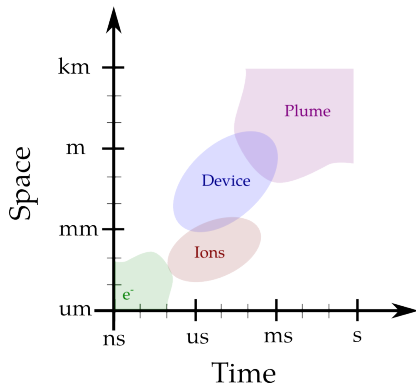
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TIME AND SPACE SCALES

Inherently Multiscale:

- $\mathcal{O}(10^9)$ -Space and Time
- Naive 3D-Spatial Scales $\rightarrow \text{Cost}^3$

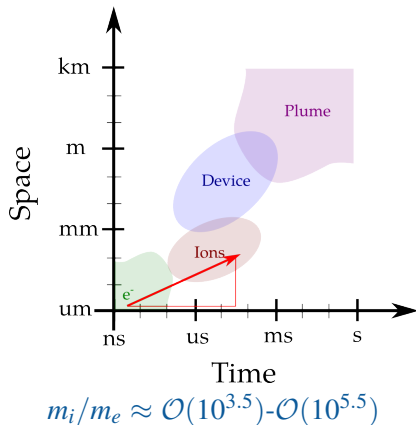




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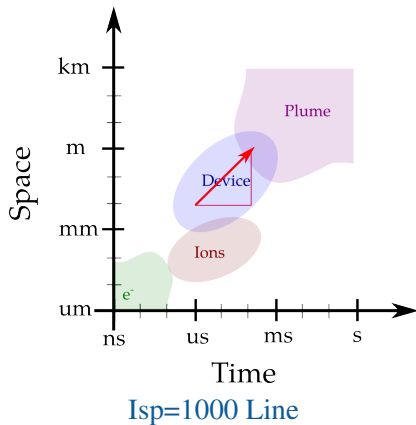




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- Isp & Device Scale \rightarrow Transit Time

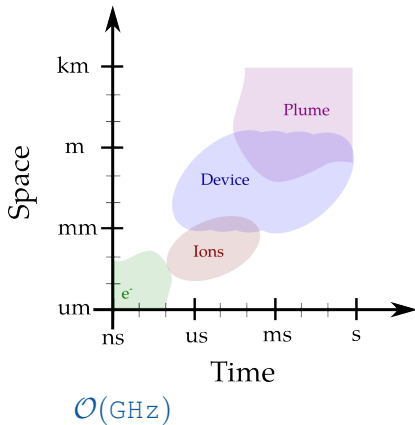




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- Isp & Device Scale \rightarrow Transit Time
- Pulsed: Spread Right at Scale & Isp

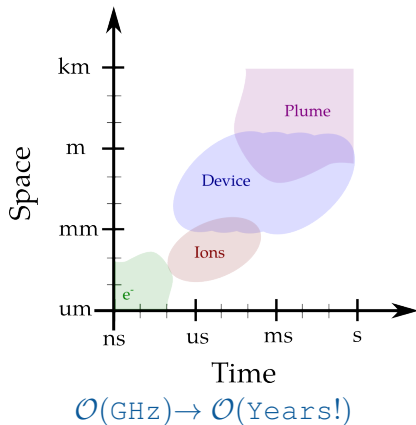




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- Spacecraft Lifetime Performance?



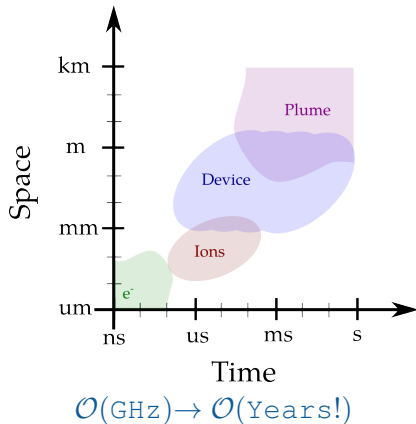


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Modeling Must Exploit Scale Separation

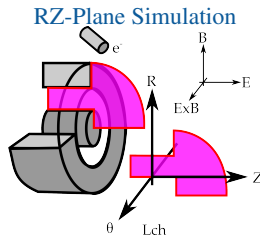




HALL EFFECT THRUSTERS

Hall Effect Thrusters (HET):

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- Named for Hall Current in θ
- Uses Quasi-1D Electron Fluid Solve
- Ohm's Law \rightarrow No e^- -momentum



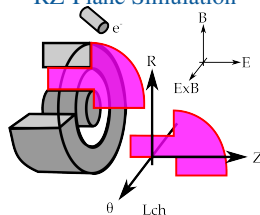


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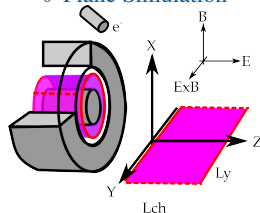
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RZ-Plane Simulation



θ -Plane Simulation



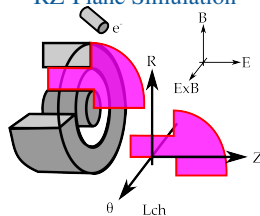


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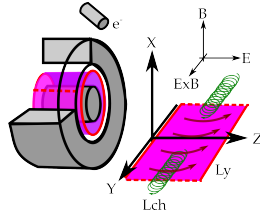
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- Results in Hall Current (Namesake)

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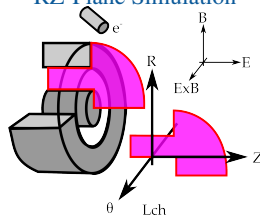


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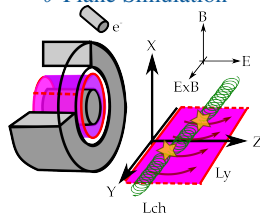
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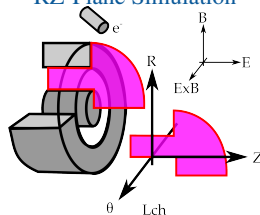


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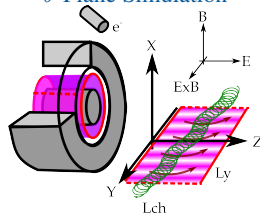
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Model Sensitive to Mobility:

- Classical Mobility Insufficient in Near Plume

From LaFleur, Phys Plasmas 23, 053503 (2016)

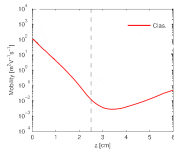


FIG. 3. Electron cross-field mobility as a function of axial position within the PPS³1350 thruster. The solid black line is an empirical mobility which is needed in fluid simulations in order to get agreement with experiment,¹⁶ the red line is the mobility based on classical diffusion across a magnetic field, while the open blue triangles show the mobility due to the saturated mobility-enhanced electron-ion friction force. The vertical dashed line indicates the thruster exit plane.

(Thruster Simulation will not “Light”)



Model Sensitive to Mobility:

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- Enhanced Empirical “Bohm” (1/B) Mobility
- Also Needs Coefficient by “Zone” (i.e. Anode/Channel/Plume Regions)

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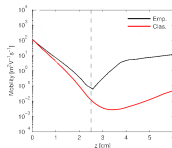
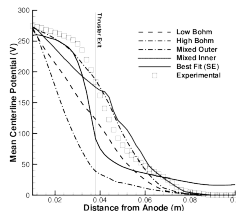


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Impact of Mobility on Discharge Profile



(Koo, PhD Dissertation)



IMPACT OF MOBILITY

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- Behavior Sensitive to Plume Coefficient

Critical if Operating Near Mode Change

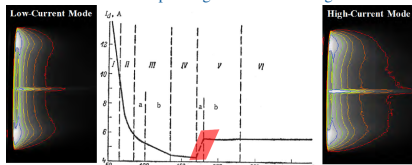
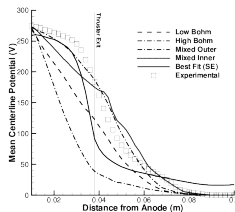


Fig. 1 Discharge current as function of magnetic field with constant discharge voltage showing operational regimes defined by Tiltman [2]
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Need Physics Based Model

Critical if Operating Near Mode Change

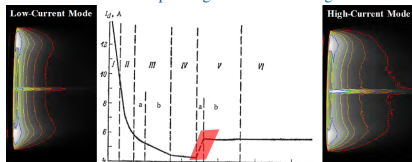
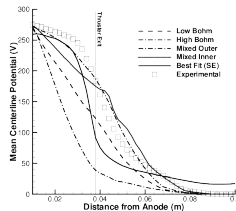


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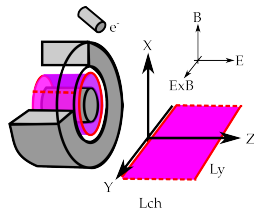


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Current Driven Instability:

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 - Anomalous e-Transport? (LaFleur)
 - Focus of 1D/2D Full-PIC Studies

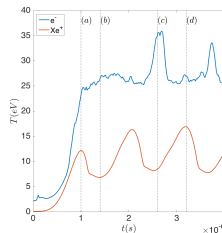
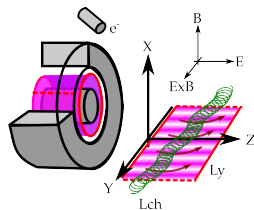




INSTABILITY AND TRANSPORT

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- $\langle \Delta n_e, \Delta E_\theta \rangle \rightarrow$ Axial Transport
Initial Exponential Growth Saturates



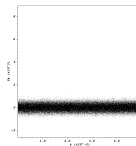


INSTABILITY AND TRANSPORT

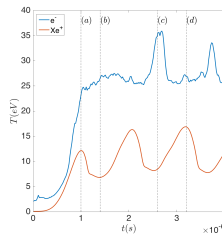
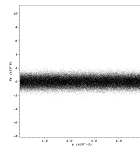
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Xe-VDF



e-VDF

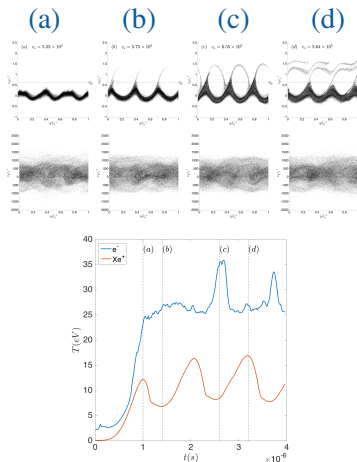




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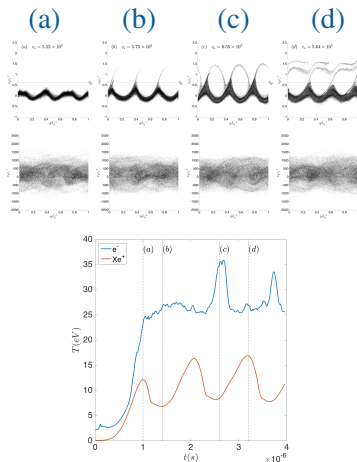




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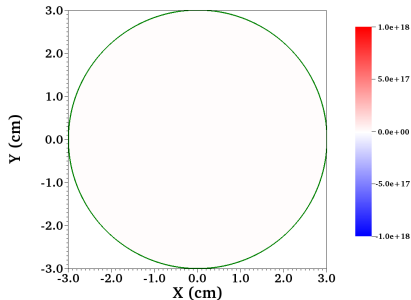




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- $j_\theta \rightarrow$ FRC Spiral Charge Separation?

Charge Separation

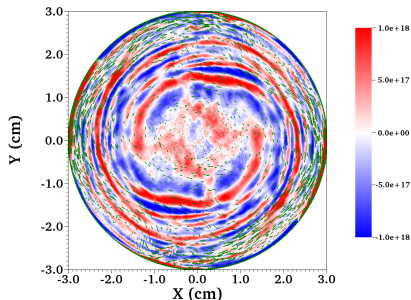




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- Onset of Plasma Turbulence?

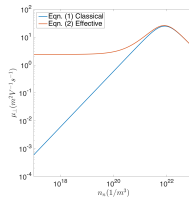
Charge Separation





Impact of Ion-Acoustic Instability:

- Instability \rightarrow Extra Mobility





Impact of Ion-Acoustic Instability:

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- LaFleur Results Promising

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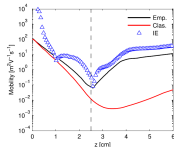
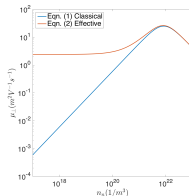


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(A Posteriori Mobility)

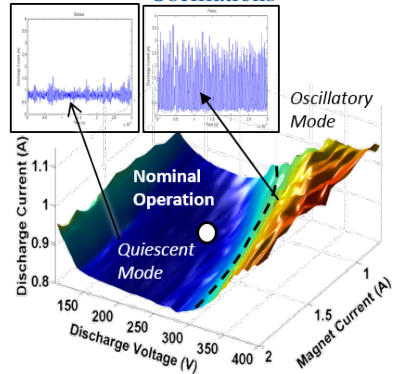




Impact of Ion-Acoustic Instability:

- Instability \rightarrow Extra Mobility
- LaFleur Results Promising
- Need High-Dim Validation

Experimental Discharge Current Oscillations



I-V-B Plot of Hall Thruster Operation

Brown, EP TEMPEST Program Review, 2015



Impact of Ion-Acoustic Instability:

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- Need High-Dim Validation
- Model via Lagged Lotka-Volterra?

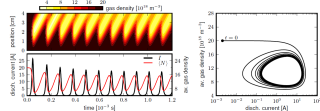


Figure 2. Convergence of the one-dimensional model (ii) to a limit cycle with a low-current, gas-filled startup. The phase portrait shows the coupled evolution of the discharge current and of the average gas density within the domain. The parameters are $\ell = 2\text{ cm}$, $V = 20\text{ cm}^3$, $Q_0 = 5 \times 10^3\text{ m}^3\text{ s}^{-1}$, $I = 4\text{ A}$, $\gamma = 4 \times 10^5\text{ s}^{-1}\text{ A}^{-1}$, $\beta_0 = 2 \times 10^5\text{ m}^3\text{ s}^{-1}$. The current profiles $i(x)$ and $\bar{n}(x)$ are shown on Fig. 1.

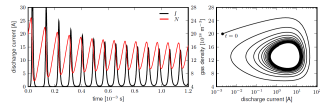


Figure 3. Convergence of the modified predator-prey model (ii) to a limit cycle with a low-current, gas-filled startup. The parameters are defined from those of Fig. 2, assuming that the length of the ionization region is given approximately by half of the discharge column length $\ell = \ell/2 = 1\text{ cm}$, $Q_0 = 5 \times 10^3\text{ m}^3\text{ s}^{-1}$, $I = 4\text{ A}$, $\gamma = 3.5 \times 10^5\text{ s}^{-1}\text{ A}^{-1}$, $\beta_0 = \beta$, $\beta_0/\beta = 3.5 \times 10^{-10}\text{ m}^3\text{ s}^{-1}$, $\tau = 2.2 \times 10^{-7}\text{ s}$.

$$\frac{dI}{dt} = \beta I(N - \bar{N}),$$

$$\frac{dN}{dt} = -\gamma IN + \frac{Q_0}{L} \exp \left[-\gamma \int_{t-\tau}^t I dt \right].$$

Barral & Peradzyński, IEPC-2009-070



Impact of Ion-Acoustic Instability:

- Instability → Extra Mobility
- LaFleur Results Promising
- Need High-Dim Validation
- Model via Lagged Lotka-Volterra?
- Model Captures a Bifurcation

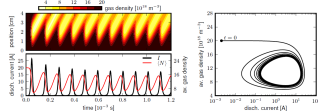


Figure 2. Convergence of the one-dimensional model (ii) to a limit cycle with a low-current, gas-filled startup. The phase portrait shows the coupled evolution of the discharge current and of the average gas density within the domain. The parameters are $\tau = 0$ ms, $V = 20$ kV, $Q_0 = 5 \times 10^3$ m³s⁻¹, $I = 4$ A, $\gamma_0 = 4 \times 10^3$ s⁻¹A⁻¹, $\beta_0 = 2 \times 10^3$ s⁻¹. The ionon profiles $\rho(x)$ and $\psi(x)$ are shown on Fig. 1.

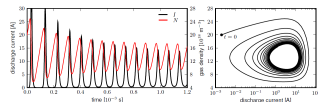


Figure 3. Convergence of the modified predator-prey model (ii) to a limit cycle with a low-current, gas-filled startup. The parameters are defined from those of Fig. 2, assuming that the length of the ionon cycle is given approximately by half of the discharge electron length, $L = l/2 = 0$ cm, $Q_0 = 5 \times 10^3$ m³s⁻¹, $I = 4$ A, $\gamma = 3.5 \times 10^3$ s⁻¹A⁻¹, $\beta = 5$ W/m³, $\beta_0 = 3.5 \times 10^3$ s⁻¹, $\tau = 2.2 \times 10^{-7}$ s.

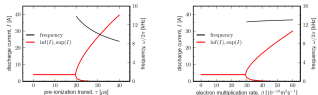


Figure 4. Bifurcation diagrams for the modified predator-prey model (ii). Parameters τ and β are varied around the nominal point investigated in Fig. 5.

$$\frac{dI}{dt} = \beta I(N - \bar{N}),$$

$$\frac{dN}{dt} = -\gamma IN + \frac{Q_0}{L} \exp \left[-\gamma \int_{t-\tau}^t I dt \right].$$

Barral & Peradzyński, IEPC-2009-070



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How can we Compare
Model, Experiment, and Simulation?

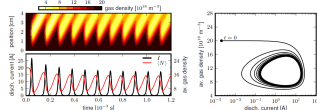


Figure 2. Convergence of the one-dimensional model (ii) to a limit cycle with a low-current, gas-filled startup. The phase portrait shows the coupled evolution of the discharge current and of the average gas density within the domain. The parameters are $\epsilon = 0.05$, $V = 20 \text{ kV}$, $Q_0 = 5 \times 10^3 \text{ m}^3 \text{ s}^{-1}$, $I = 4 \text{ A}$, $\gamma_0 = 4 \times 10^3 \text{ s}^{-1} \text{ A}^{-1}$, $\beta_0 = 2 \times 10^3 \text{ s}^{-1}$. The ionon profiles $\rho(x)$ and $\psi(x)$ are shown on Fig. 1.

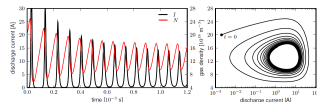


Figure 3. Convergence of the modified protogrey model (ii) to a limit cycle with a low-current, gas-filled startup. The parameters are defined from those of Fig. 2, assuming that the length of the ionon cycle is given approximately by half of the discharge relaxation length, $L = l/2 = 0.05 \text{ m}$, $Q_0 = 5 \times 10^3 \text{ m}^3 \text{ s}^{-1}$, $I = 4 \text{ A}$, $\gamma = 3.5 \times 10^3 \text{ s}^{-1} \text{ A}^{-1}$, $\beta = 5 \text{ W/Hz}$, $\beta_0 = 3.5 \times 10^3 \text{ s}^{-1} \text{ m}^3 \text{ s}^{-1}$, $\tau = 2.2 \times 10^{-7} \text{ s}$.

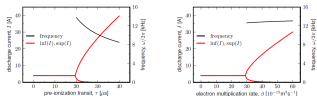


Figure 4. Bifurcation diagrams for the modified protogrey model (ii). Parameters τ and β are varied around the nominal point investigated in Fig. 5.

$$\frac{dI}{dt} = \beta I(N - \bar{N}),$$

$$\frac{dN}{dt} = -\gamma N + \frac{Q_0}{L} \exp \left[-\gamma \int_{t-\tau}^t I dt \right].$$

Barral & Peradzyński, IEPC-2009-070



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How can we Compare
Model, Experiment, and Simulation?
With only **I(t)** Accessible Experimentally?

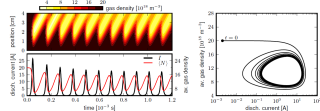


Figure 2. Convergence of the one-dimensional model (in a finite cycle with a low-current, gas-filled startup). The phase portrait shows the coupled evolution of the discharge current and of the average gas density within the domain. The parameters are $\tau = 10\text{ ns}$, $V = 20\text{ kV}$, $Q_0 = 5 \times 10^3\text{ m}^{-3}\text{ s}^{-1}$, $I = 4\text{ A}$, $\gamma_0 = 4 \times 10^3\text{ s}^{-1}\text{ A}^{-1}$, $\beta_0 = 2 \times 10^3\text{ s}^{-1}\text{ m}^{-3}$. The nominal profiles $\rho(x)$ and $\psi(x)$ are shown on Fig. 1.

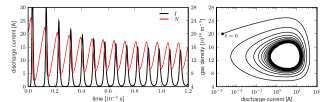


Figure 3. Convergence of the modified protodecay model (in a finite cycle with a low-current, gas-filled startup). The parameters are defined from those of Fig. 2, assuming that the length of the relaxation cycle is given approximately by half of the discharge relaxation length, $L = l/2 = 10\text{ cm}$, $Q_0 = 5 \times 10^3\text{ m}^{-3}\text{ s}^{-1}$, $I = 4\text{ A}$, $\gamma = 3.5 \times 10^3\text{ s}^{-1}\text{ A}^{-1}$, $\beta = 5\text{ W/m}^3$, $\beta_0 = 3.5 \times 10^3\text{ m}^{-3}\text{ s}^{-1}$, $\tau = 2.2 \times 10^{-7}\text{ s}$.

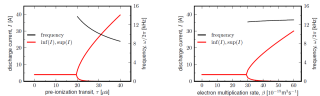


Figure 4. Bifurcation diagrams for the modified protodecay model (9). Parameters τ and β are varied around the nominal point investigated in Fig. 5.

$$\frac{dI}{dt} = \beta I(N - \bar{N}),$$

$$\frac{dN}{dt} = -\gamma IN + \frac{Q_0}{L} \exp \left[-\gamma \int_{t-\tau}^t I dt \right].$$

Barral & Peradzyński, IEPC-2009-070



State Space Reconstruction: Time Series and Dynamic Systems

A supplemental simulation and animation for
"Detecting Causality in Complex Ecosystems"

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animation by: Peter Sugihara, Hao Ye, and George Sugihara

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Same Could be Done for Hall Thruster Simulation



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Same Could be Done for Hall Thruster Simulation
But Experiment has Limited Access?



State Space Reconstruction: Takens' Theorem and Shadow Manifolds

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Shadow Maifold from only $I(t)$?



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Shadow Maifold from only $I(t)$?
Useful Tuning Model/Simulation Parameters?



State Space Reconstruction: Convergent Cross Mapping

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Useful Detecting Causal Relationships.



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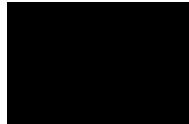
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Useful Detecting Causal Relationships.
Potential Tool for Self-Similarity/Causality in Turbulence?



DMD → Turbulence:

- Low Re Vortex Shedding



Von Karman Vortex Street

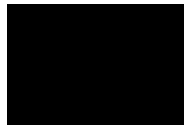
http://commons.wikimedia.org/wiki/File:Karman_Vortex_Street_Off_Cylinder.ogg



TURBULENCE FROM DYNAMIC SYSTEM

DMD → Turbulence:

- Low Re Vortex Shedding
- Sparse Dynamic System via DMD



Von Karman Vortex Street

http://commons.wikimedia.org/wiki/File:Karman_Vortex_Street_Off_Cylinder.orgv

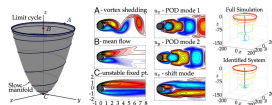


Fig. 2. Example of high-dimensional dynamical system from fluid dynamics. The vortex shedding past a cylinder is a prototypical example that is used for flow control, with relevance to many applications, including drag reduction behind vehicles. The vortex shedding is the result of a global bifurcation. However, because the Navier-Stokes equations have quadratic nonlinearity, it is necessary to use a mean-field model with a separation of timescales, where a fast mean-field deformation is slow to the slow vortex shedding dynamics. The parabolic slow manifold is shown (left), with the unstable fixed point (C), mean flow (B), and vortex shedding (A). A POD basis and shift mode are used to reduce the dimension of the problem (middle right). The identified dynamics closely match the true trajectory in POD coordinates, and most importantly, they capture the quadratic nonlinearity and timescales associated with the mean-field model.

Brunton, Proctor, and Kutz, 3932-3937, PNAS, 4/12/16, v.113, no. 15



DMD \rightarrow Turbulence:

- Low Re Vortex Shedding
- Sparse Dynamic System via DMD
- What if Flow is Non-Sparse? (Turbulent)

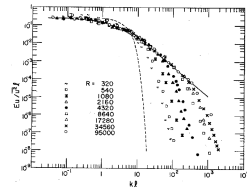


FIG. 6. Dependence of the frequency spectra $E_u / (u^2 l)$ on Reynolds number in the center of the wake.

Ubeeroi and Freymuth, Phys. Fluids 12, 1359 (1969).



DMD \rightarrow Turbulence:

- Low Re Vortex Shedding
- Sparse Dynamic System via DMD
- What if Flow is Non-Sparse? (Turbulent)
- Inertial Range Turbulence really Self-Similar?
- Universal if Agnostic to Large and Small k

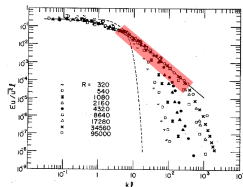


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TURBULENCE FROM DYNAMIC SYSTEM

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- Bandpass Invariant up to Similarity?
(Filter Local Space/Time: Gabor? Wavelets?)

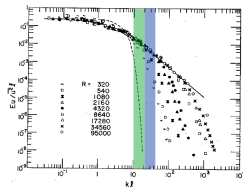
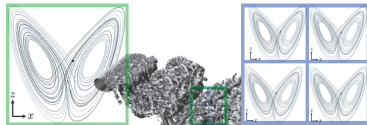


Fig. 6. Dependence of the frequency spectra $E_u/\sqrt{u}^3 l$ on Reynolds number in the center of the wake.

Ubroeri and Freymuth, Phys. Fluids 12, 1359 (1969).



Piomelli, Phil. Trans. Royal Soc., A 372, 2014.

https://commons.wikimedia.org/wiki/File:A_Trajectory_Through_Phase_Space_in_a_Lorenz_Attractor.gif



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(Filter Local Space/Time: Gabor? Wavelets?)
- Inter-Band Dynamics Causal via CCM?

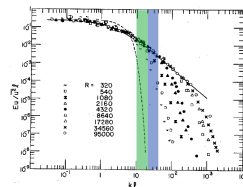
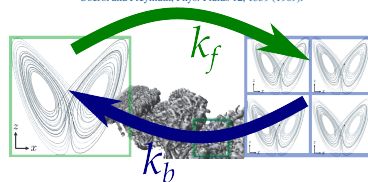


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(Filter Local Space/Time: Gabor? Wavelets?)
- Inter-Band Dynamics Causal via CCM?
- CCM → K_j in Mori-Zwanzig LES?

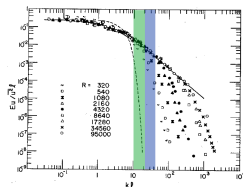
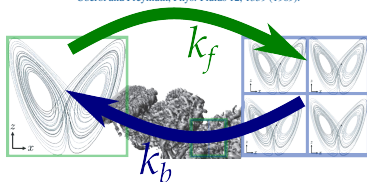


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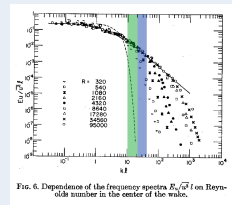
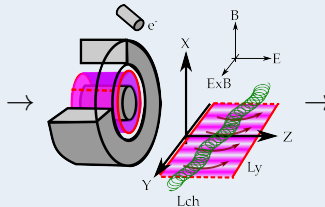
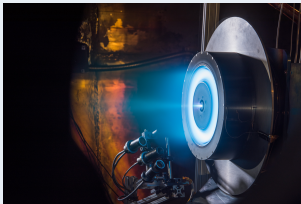
$$w_j^{(0)} = \mathcal{P} \int_0^t K_j(\bar{u}(t-s), s) ds$$

Parish and Duraisamy, AIAA Expo, 6/16



SUMMARY & CONCLUSION

Thank You



Work Supported through AFOSR Task 17RQCOR465 (PM: Birkan)

Questions?